p-ADIC STRING THEORY

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- Introduction
- P-Adic numbers, adeles and their functions
- P-Adic strings
- Concluding remarks

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1. Introduction

p-adic mathematical physics (1987)

- p-adic and adelic string theory
- p-adic and adelic quantum mechanics and quantum field theory
- p-adic and adelic gravity and (quantum) cosmology
- *p*-adic and adelic space-time structure at the Planck scale
- p-adic and adelic dynamical systems
- p-adic stochastic processes
- *p*-adic aspects of information theory
- p-adic wavelets
- spin glasses
- protein dynamics
- p-adic structure of the genetic code

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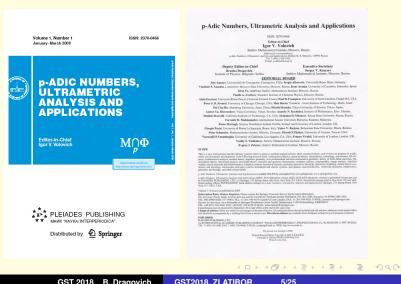
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Some review references

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- V.S. Vladimirov, I.V. Volovich and E.I. Zelenov, *p-Adic Analysis* and Mathematical Physics (World Scientific, Singapore, 1994).
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1. Introduction: international journal on p-adics



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1. Introduction: international conferences

 Moscow 2003, Belgrade 2005, Moscow 2007, Grodno 2009, Bielefeld 2013, Belgrade 2015, Mexico 2017, Portugal 2019

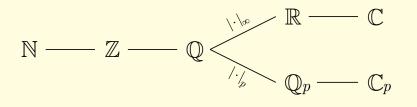


2nd International Conference on p-Adic Mathematical Physics Belgrade, 2005

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- Ostrowski theorem: Any nontrivial norm on Q is equivalent to usual absolute value or to *p*-adic norm, where *p* is any prime number (*p* = 2, 3, 5, 7, 11, ...).
- *p*-adic norm of $x \in \mathbb{Q}$: $|x|_p = |p^{\nu} \frac{a}{b}|_p = p^{-\nu}, \quad \nu \in \mathbb{Z}.$
- \mathbb{Q} is dense in \mathbb{Q}_p and \mathbb{R} .
- Completion of Q with respect to *p*-adic distance gives the field Q_p of *p*-adic numbers, in analogous way to construction of the field ℝ of real numbers.



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Examples *p*-adic norms of $x = \frac{21}{20} = \frac{3 \cdot 7}{2^2 \cdot 5}$ • $|\frac{21}{20}|_2 = 4$ • $|\frac{21}{20}|_3 = \frac{1}{3}$ • $|\frac{21}{20}|_5 = 5$ • $|\frac{21}{20}|_7 = \frac{1}{7}$ • $|\frac{21}{20}|_p = 1$ if $p \ge 11$. *p*-adic expansion of -1

$$-1 = (p-1) + (p-1)p + (p-1)p^2 + (p-1)p^3 + \dots$$

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- *p*-Adic numbers invented by Kurt Hensel (1861-1941) in 1897.
- Any *p*-adic number (*x* ∈ Q_{*p*}) has a unique canonical representation

$$x = p^{\nu} \sum_{n=0}^{+\infty} x_n p^n, \ \nu \in \mathbb{Z}, \ x_n \in \{0, 1, \cdots, p-1\}$$

$$|x_n p^{(\nu+n)}|_p = p^{-(\nu+n)}, \quad |x+y|_p \le max\{|x|_p, |y|_p\}$$

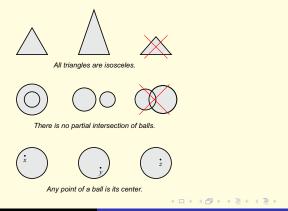


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- *p*-adic distances d_p(x, y) = |x − y|_p are non-Archimedean (ultrametric): d_p(x, y) ≤ max{d_p(x, z), d_p(z, y)}
- open ball of radius *r* and centre *a*:
- closed ball of radius r and centre a:
- balls are simultaneously closed and open



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Real and *p*-adic numbers are unified by adeles. An adele *α* is an infinite sequence

$$\alpha = (\alpha_{\infty}, \alpha_{2}, \alpha_{3}, \cdots, \alpha_{p}, \cdots), \quad \alpha_{\infty} \in \mathbb{R}, \ \alpha_{p} \in \mathbb{Q}_{p}$$

where for all but a finite set *S* of primes *p* one has that $\alpha_p \in \mathbb{Z}_p = \{x \in \mathbb{Q}_p : |x|_p \le 1\}$, i.e. *p*-adic integers. • Space of adeles

$$\mathbb{A} = \bigcup_{S} \mathcal{A}(S), \quad \mathcal{A}(S) = \mathbb{R} \times \prod_{p \in S} \mathbb{Q}_p \times \prod_{p \notin S} \mathbb{Z}_p.$$



C. Chevalley (1909 - 1984)

A. Weil (1906 - 1998

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Recall standard theoretical models

classical models : $\mathbb{R} \to \mathbb{R}$

quantum models : $\mathbb{R} \to \mathbb{C} \to \mathbb{R}$

p-Adic modeling

classical models : $\mathbb{Q}_{p} \to \mathbb{R}$

quantum models : $\mathbb{Q}_p \to \mathbb{C} \to \mathbb{R}$

 p-Adic side of some phenomena may emerge at a deeper level.

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Some connections of *p*-adic and real norms of rational numbers

$$|x|_{\infty} imes \prod_{oldsymbol{
ho}\in\mathbb{P}}|x|_{oldsymbol{
ho}}=1\ , \ ext{if} \ x\in\mathbb{Q}^{ imes}$$

$$\chi_{\infty}(x) imes \prod_{p \in \mathbb{P}} \chi_{p}(x) = 1, ext{ if } x \in \mathbb{Q}$$

 $\chi_{\infty}(x) = \exp(-2\pi i x), \quad \chi_{p}(x) = \exp 2\pi i \{x\}_{p}$

$$egin{aligned} &\prod_{oldsymbol{p}\in\mathbb{A}}\Omega(|x|_{oldsymbol{p}}) = egin{cases} &1, &x\in\mathbb{Z},\ &0, &x\in\mathbb{Q}\setminus\mathbb{Z} \end{aligned} \ &\Omega(|x|_{oldsymbol{p}}) = egin{cases} &1, &|x|_{oldsymbol{p}}\leq 1,\ &0, &|x|_{oldsymbol{p}}>1 \end{aligned}$$

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3. *p*-Adic strings

Volovich, Freund, Witten, Vladimirov, Arefeva, B. D., ... String amplitudes (1987):

standard crossing symmetric Veneziano amplitude

$$egin{aligned} \mathcal{A}_{\infty}(a,b) &= g_{\infty}^2 \, \int_{\mathbb{R}} |x|_{\infty}^{a-1} \, |1-x|_{\infty}^{b-1} \, d_{\infty} x \ &= g_{\infty}^2 \, rac{\zeta(1-a)}{\zeta(a)} \, rac{\zeta(1-b)}{\zeta(b)} \, rac{\zeta(1-c)}{\zeta(c)} \end{aligned}$$

p-adic crossing symmetric Veneziano amplitude

$$egin{aligned} &A_{p}(a,b) = g_{p}^{2} \int_{\mathbb{Q}_{p}} |x|_{p}^{a-1} \, |1-x|_{p}^{b-1} \, d_{p}x \ &= g_{p}^{2} \, rac{1-p^{a-1}}{1-p^{-a}} \, rac{1-p^{b-1}}{1-p^{-b}} \, rac{1-p^{c-1}}{1-p^{-c}} \end{aligned}$$

where a = -s/2 - 1 and $a, b, c \in \mathbb{C}$ and a + b + c = 1.

3. *p*-Adic strings

- *p*-Adic strings are strings which world-sheet is *p*-adic.
- Freund-Witten (adelic) product formula for ordinary and p-adic strings

$$A(a,b) = A_{\infty}(a,b) \prod_{\rho} A_{\rho}(a,b) = g_{\infty}^2 \prod_{\rho} g_{\rho}^2 = const.$$

- Ordinary and p-adic strings are at the equal footing
- Amplitude for real string A_∞(a, b), which is a special function, can be presented as product of inverse p-adic amplitudes, which are elementary functions.

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3. *p*-Adic strings – Effective Lagrangian

Freund, Witten, Frampton, Okada, ...

- There is an effective field description of scalar open and closed *p*-adic strings. The corresponding Lagrangians are very simple and exact. They describe not only four-point scattering amplitudes but also all higher ones at the tree-level.
- The exact tree-level Lagrangian for effective scalar field φ which describes open p-adic string tachyon is

$$\mathcal{L}_{p} = \frac{m_{p}^{D}}{g_{p}^{2}} \frac{p^{2}}{p-1} \left[-\frac{1}{2} \varphi p^{-\frac{\Box}{2m_{p}^{2}}} \varphi + \frac{1}{p+1} \varphi^{p+1} \right]$$

where $\Box = -\partial_t^2 + \nabla^2$ is the *D*-dimensional d'Alembertian and metric with signature (- + ... +).

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3. *p*-Adic strings

$$\mathcal{L}_{p} = rac{m_{p}^{D}}{g_{p}^{2}} rac{p^{2}}{p-1} \Big[-rac{1}{2} \, arphi \, p^{-rac{\Box}{2m_{p}^{2}}} \, arphi + rac{1}{p+1} \, arphi^{p+1} \Big]$$

• Lagrangian with *p*-adic world-sheet B.D.

$$\begin{split} \mathcal{L}_{p} = & \frac{m^{D}}{g^{2}} \frac{p^{2}}{p-1} \Big[\frac{1}{2} \varphi \int_{\mathbb{R}} \Big(\int_{\mathbb{Q}_{p} \setminus \mathbb{Z}_{p}} \chi_{p}(u) |u|_{p}^{\frac{k^{2}}{2m^{2}}} du \Big) \tilde{\varphi}(k) \, \chi_{\infty}(kx) \, d^{4}k \\ &+ \frac{1}{p+1} \, \varphi^{p+1} \Big] \\ &\int_{|x|_{p} > 1} \chi_{p}(u) |u|_{p}^{s} du = -p^{s} \end{split}$$

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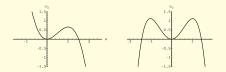
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• The corresponding potential $\mathcal{V}(\varphi)$ for this Lagrangian is $\mathcal{V}_{p}(\varphi) = -\mathcal{L}_{p}(\Box = 0)$, which the explicit form is

$$\mathcal{V}_{p}(\varphi) = rac{m^{D}}{g^{2}} \left[rac{1}{2} rac{p^{2}}{p-1} \varphi^{2} - rac{p^{2}}{p^{2}-1} \varphi^{p+1}
ight]$$

It has local minimum $V_p(0) = 0$. If $p \neq 2$ there are two local maxima at $\varphi = \pm 1$ and there is one local maximum $\varphi = +1$ when p = 2.

The 2-adic string potential V₂(φ) (on the left) and 3-adic potential V₃(φ) (on the right) of the figure.



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3. *p*-Adic strings – equation of motion

The equation of motion is

$$\varphi^{-\frac{\Box}{2m_p^2}}\varphi=\varphi^p$$

- It has trivial solutions φ = 0 and φ = 1, and φ = −1 for p ≠ 2.
- There are also solutions in any direction x_i and time t

$$\varphi(x_i) = p^{\frac{1}{2(p-1)}} \exp\left(-\frac{p-1}{2p\ln p}m^2 x_i^2\right)$$

$$\varphi(t) = p^{\frac{1}{2(p-1)}} \exp\left(\frac{p-1}{2p\ln p}m^2t^2\right)$$

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3. *p*-Adic strings

There is also D-dimensional solution

$$\varphi(x) = p^{\frac{D}{2(p-1)}} \exp\left(-\frac{p-1}{2p\ln p}m^2x^2\right), \quad x^2 = -t^2 + \sum_{i=1}^{D-1}x_i^2$$

The above solutions can be obtained using identity

$$e^{A\partial_t^2} e^{Bt^2} = rac{1}{\sqrt{1-4AB}} e^{rac{Bt^2}{1-4AB}}, \quad 1-4AB > 0.$$

p-Adic realization of tachyon condensation

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3. *p*-Adic strings – Lagrangian for *p*-adic string sector

B.D. Consider

$$L = \sum_{n=1}^{+\infty} C_n \mathcal{L}_n = m^D \sum_{n=1}^{+\infty} \frac{C_n}{g_n^2} \frac{n^2}{n-1} \left[-\frac{1}{2} \phi \, n^{-\frac{\Box}{2m^2}} \phi + \frac{1}{n+1} \, \phi^{n+1} \right]$$

where $\frac{C_n}{g_n^2} \frac{n^2}{n-1} = (-1)^{n-1}$. Take into account

$$\sum_{n=1}^{+\infty} (-1)^{n-1} \frac{1}{n^s} = (1-2^{1-s}) \zeta(s), \quad s = \sigma + i\tau, \quad \sigma > 0$$

which has analytic continuation to the entire complex *s* plane without singularities.

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3. *p*-Adic strings – Lagrangian for *p*-adic string sector

• The corresponding Lagrangian is

$$L = m^{D} \left[-\frac{1}{2} \phi \left(1 - 2^{1 - \frac{\Box}{2m^{2}}} \right) \zeta \left(\frac{\Box}{2m^{2}} \right) \phi + \phi - \frac{1}{2} \log(1 + \phi)^{2} \right].$$

The potential is

$$V(\phi) = -L(\Box = 0) = m^D \Big[rac{1}{4} \phi^2 - \phi + rac{1}{2} \log(1 + \phi)^2 \Big],$$

which has one local maximum V(0) = 0 and one local minimum at $\phi = 1$. It is singular at $\phi = -1$.

The equation of motion is

$$\left(1-2^{1-\frac{\Box}{2m^2}}\right)\zeta\left(\frac{\Box}{2m^2}\right)\phi=\frac{\phi}{1+\phi},\quad\phi=0.$$

4. Concluding remarks

• Taking limit $p = 1 + \varepsilon \rightarrow 1$ in effective Lagrangian, one obtains

$$L_1(\varphi) = \frac{m^D}{g^2} \Big[\frac{1}{2} \varphi \frac{\Box}{m^2} \varphi + \frac{\varphi^2}{2} (\ln \varphi^2 - 1) \Big]$$

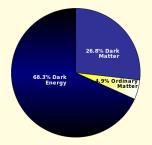
which corresponds to the ordinary bosonic string in the boundary string field theory.

- p-Adic tachyon condensation
- *p*-Adic inflation
- p-Adic AdS/CFT correspondence



4. Concluding remarks

- If Einstein general theory of relativity is theory of gravity for the Universe as a whole, then there is only about 5% of ordinary matter, about 27% of dark matter and about 68% of dark energy.
- Conjecture: Dark matter and dark energy have *p*-adic origin.
- Modified gravity may be nonlocal, e.g. *R* → *R* + *RF*(□)*R*.



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4. Concluding remarks

- *p*-Adic numbers are not results of measurements, but can be useful in physical models.
- *p*-Adic strings are strings with world-sheet labeled by *p*-adic numbers.
- If there exist ordinary strings then there should exist also p-adic strings.
- *p*-Adic string theory is simpler than ordinary string theory, and is useful for ordinary string theory.
- *p*-Adic strings may play significant role in description of dark side of the universe.